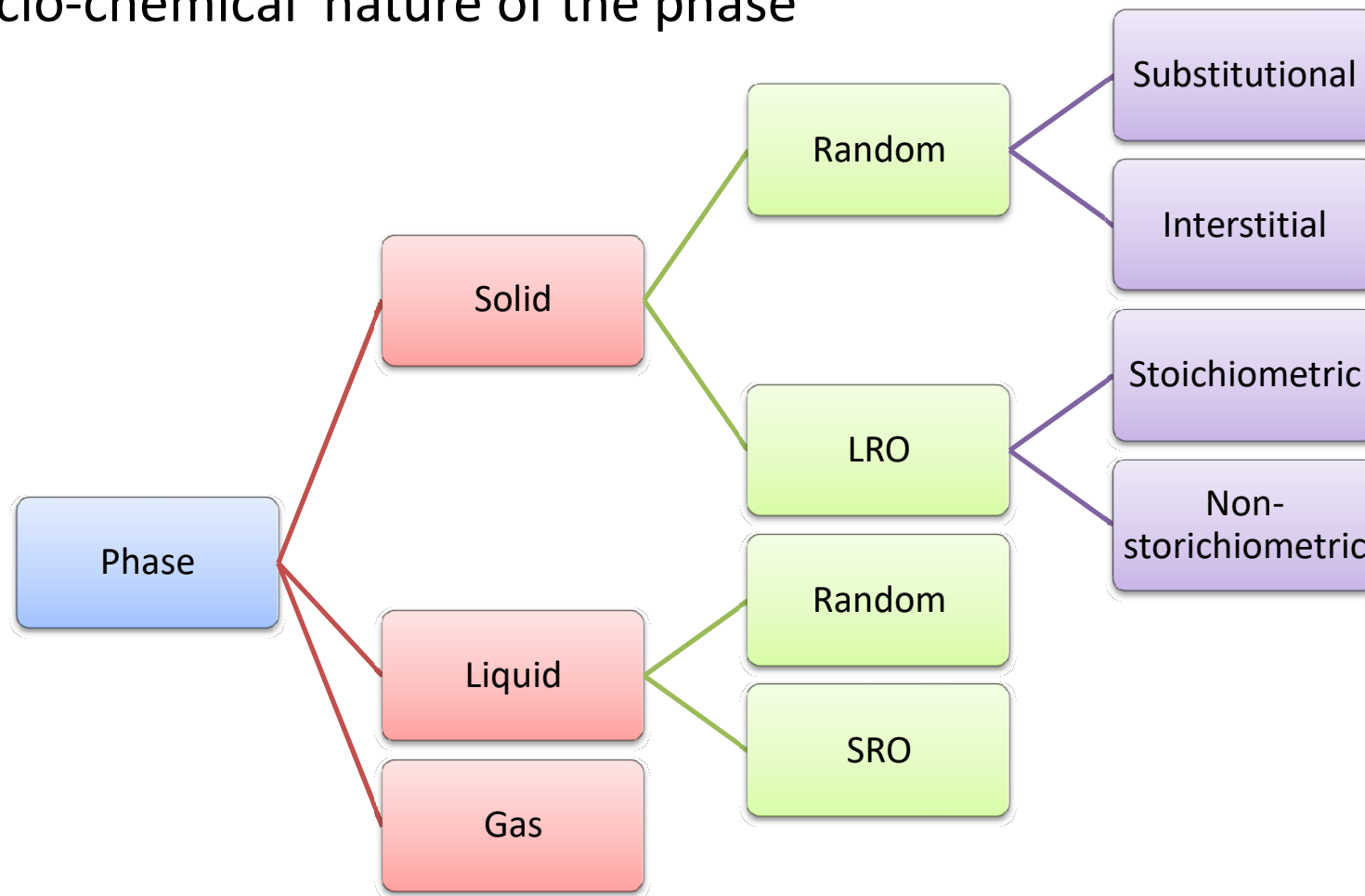


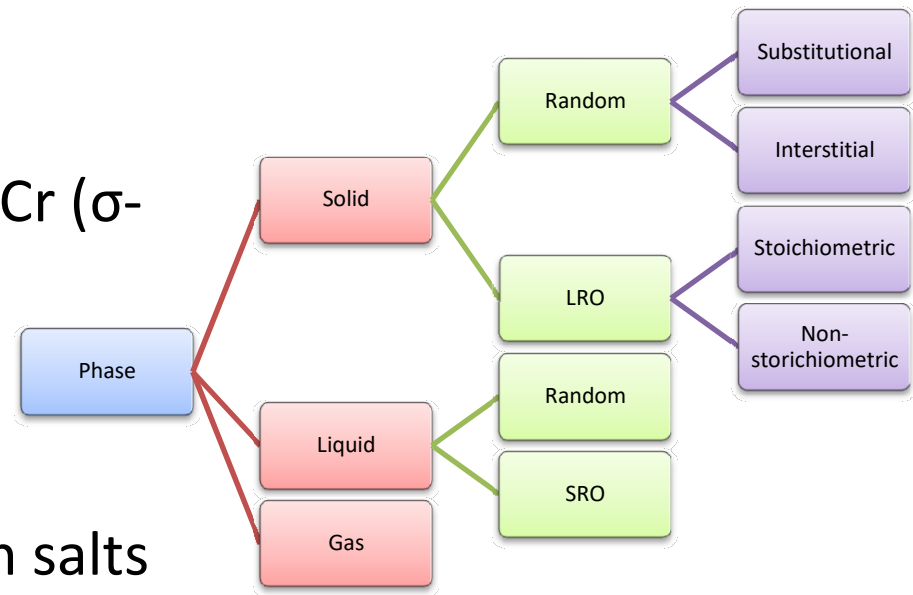
# Gibbs Energy Model

# Gibbs Energy Models

- Models for T, P and composition dependence of Gibbs energy
- Choice of a particular Gibbs energy model depends on the physico-chemical nature of the phase



- Random-Substitutional:  $\alpha$ -brass
- Random-Interstitial: C in ( $\alpha$ -Fe)
- LRO Non-stoichiometric: NiAl, FeCr ( $\sigma$ -phase), FeO (Wustite)
- LRO-Stoichiometric:  $\text{Fe}_3\text{C}$ ,  $\text{Al}_2\text{O}_3$
- Liquid-Random: molten steel
- Liquid-SRO: slags, mattes, molten salts



# Elements and Stoichiometric Phases

- G has no composition dependence,

$$G(T, p) = G^{lattice} + G^{phys} + G^{pres}$$

where  $G^{lattice}$  is the lattice (chemical) contribution

$G^{phys}$  is the physical contribution like magnetic, ordering

$G^{pres}$  is the pressure contribution

- For elements one should use G functions recommended by SGTE
- SGTE expression for  $G^{lattice}$  of an element i has the following format (Meyer-Kelly Polynomial):

$$G_i^0(T, 1bar) = H_i^{SER} + a + bT + cT \ln T + \sum_{n=2,3,-1} d_n T^n$$

Where a, b, c, ... are the model parameters

$H_i^{SER}$  serves as the reference energy term

$$G_i^0(T, 1bar) = H_i^{SER} + a + bT + cT \ln T + \sum_{n=2,3,-1} d_n T^n$$

- SER refers to Standard Element Reference, the reference state adopted by SGTE (Scientific Group Thermodata Europe)
- According to SER,

$$H_i^{SER} = H_i^{0,\phi}(298.15K, 1bar) = 0$$

$$S_i^{0,\phi}(0K, 1bar) = 0$$

$$S_i^{0,\phi}(298.15K, 1bar) = \int_0^{298.15} C_{pi}^{0,\phi} d \ln T$$

Where  $\phi$  corresponds to the phase of  $i$  that is stable at 298.15K and 1 bar.

NOTE: SGTE expressions are valid only above 298.15 K

$$G_i^0(T, 1\text{bar}) = H_i^{SER} + a + bT + cT \ln T + \sum_{n=2,3,-1} d_n T^n$$

- Following can be deduced from the SGTE expression for

 $G_i^0$ 

$$S_i^0 = -b - c(1 + \ln T) - \sum_{n=2,3,-2} d_n T^{n-1}$$

$$H_i^0 - H_i^{SER} = a - cT - \sum_{n=2,3,-2} (n-1)d_n T^n$$

$$C_{p,i}^0 = -c - \sum_{n=2,3,-2} n(n-1)d_n T^{n-1}$$

- For **stoichiometric phases (compounds)**  $G^{\text{lattice}}$  is represented as

$$G^0(T, 1\text{bar}) = \sum_i \nu_i H_i^{\text{SER}} + A + BT + CT \ln T + \sum_{n=2,3,-1} D_n T^n$$

where  $\nu_i$  are the stoichiometric coefficients

A, B, C, ... are the model parameters.

For compounds with no  $C_p$  of formation (Neumann and Kopp rule,  $\Delta_f C_p = 0$ ), one can use a simple format

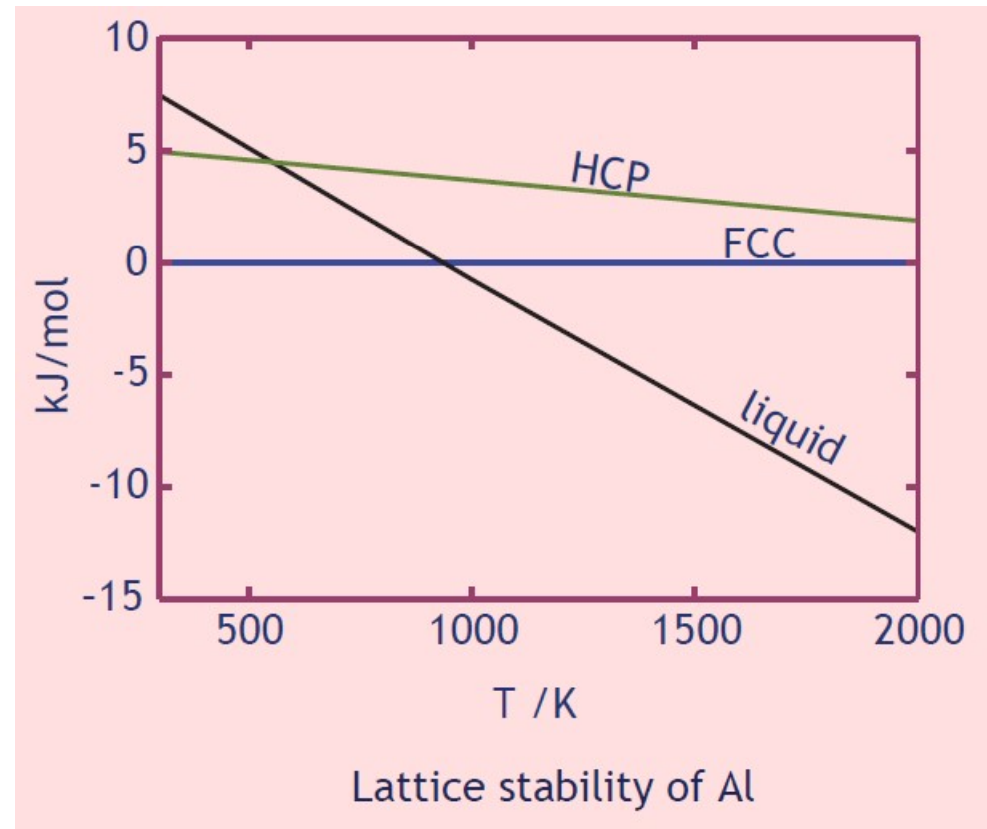
$$G^0(T, 1\text{bar}) = \sum_i \nu_i G_i^{0,ref} + \Delta_f G$$

$$\Delta_f G = h + sT$$

Where  $\Delta_f H = h$  and  $\Delta_f S = s$

# Lattice Stability

- Gibbs energy difference between two phases of an element is called lattice stability.
- Ex.:  $G_{Al}^{0,HCP} - G_{Al}^{0,FCC}$





## Example

$$G_C^{0,gr} = H_C^{SER} - 17368.441 + 170.73T - 24.3T \ln(T) \\ - 4.723 \times 10^{-4} + \frac{2562600}{T} - \frac{2.643 \times 10^8}{T^2} + \frac{1.2 \times 10^{10}}{T^3}$$

$$G^{0,cem} = 3H_{Fe}^{SER} + H_C^{SER} - 10745 + 706.04T - 120.6T \ln(T)$$

$$G^{0,AlB_2} = G_{Al}^{0,FCC} + G_B^{0,rhomb} - 85808.76 + 45.46923T$$

$$G_{Al}^{0,HCP} - G_{Al}^{0,FCC} = +5481 - 1.8T$$

- $T$  dependence of heat capacity is no longer simple when phenomena such as magnetic transitions are involved
- Accounting Gibbs energy due to such physical phenomena separately is more appropriate
- **Inden-Hillert-Jarl model** for magnetic transitions

$$G^{ma} = RT \ln(\beta + 1)g(\tau)$$

where,  $\tau = T/T_c$

$T_c$  is the critical temperature for the magnetic transition

$\beta$  is the mean magnetic moment expressed in Bohr magnetons

The function  $g(\tau)$  is

$$g(\tau) = 1 - \frac{1}{A} \left( \frac{79\tau^{-1}}{140} + \frac{474}{497} \left( \frac{1}{s} - 1 \right) \left( \frac{\tau^3}{6} + \frac{\tau^9}{135} + \frac{\tau^{15}}{600} \right) \right) \quad \tau < 1$$

$$g(\tau) = -\frac{1}{A} \left( \frac{\tau^{-5}}{10} + \frac{\tau^{-15}}{315} + \frac{\tau^{-25}}{1500} \right) \quad \tau \geq 1$$

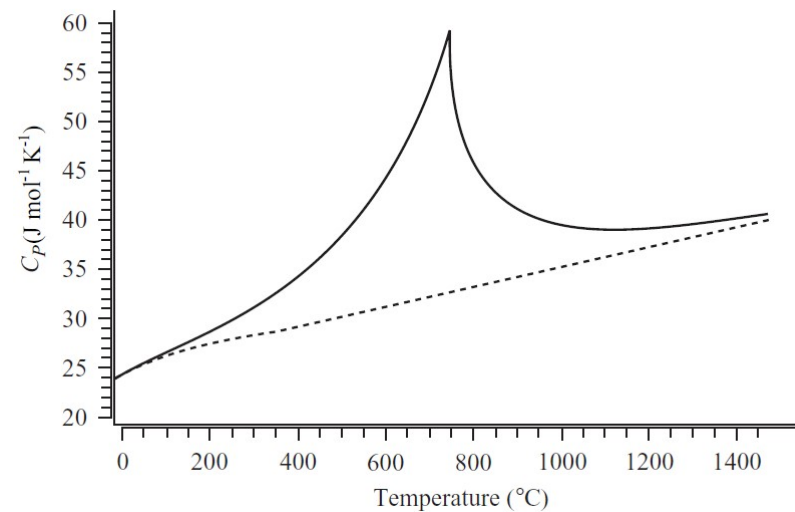
Where,

$$A = \frac{518}{1125} + \frac{11692}{15975} \left( \frac{1}{p} - 1 \right)$$

P is a structure dependent constant  $p = \frac{(s+1)}{s(z-1)}$

Z is the CN (Z=8 for BCC, Z=12 for FCC, HCP)

S (=1/2) electron spin quantum number



Heat capacity of  $\alpha$ -Fe

# Gas Phase

- Gas (Vapour) phase may consist of several species.  
Example: with components O and S we have O, O<sub>2</sub>, O<sub>3</sub>, S, S<sub>2</sub> etc. constituent species.
- It is common to treat gas phase as ideal ( $V=RT/p$ )

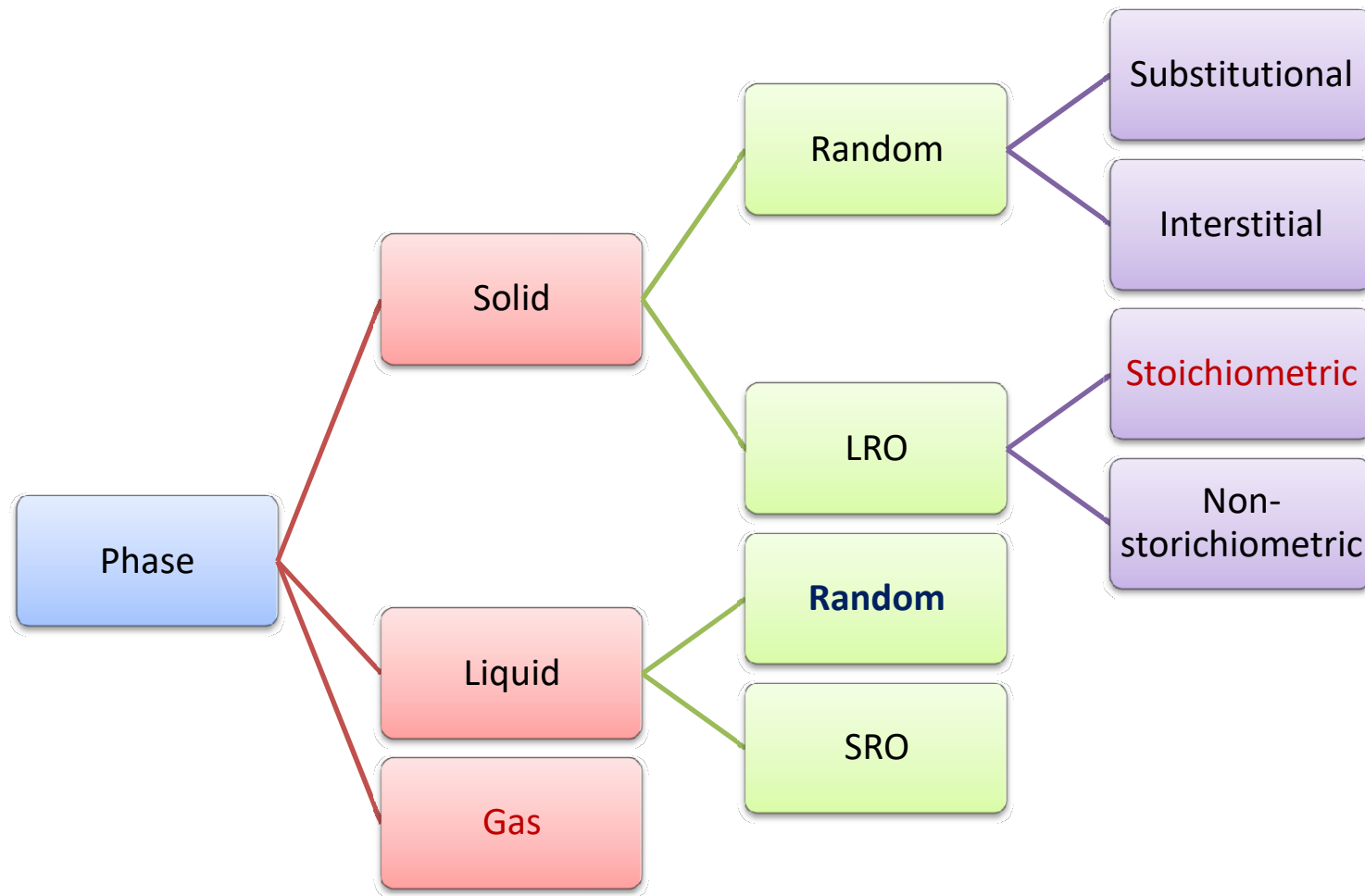
$$G = G^{ref} + G^{conf} + G^{pres}$$
$$G = G^{ref} - TS^{conf} + G^{pres}$$
$$G^{ref} = \sum_i y_i \left( G_i^{0,gas} - \sum_j b_{ij} H_i^{SER} \right)$$
$$-TS^{conf} = RT \sum_i y_i \ln y_i$$
$$G^{pres} = RT \int_{p^0}^p \frac{dp}{p} = RT \ln \left( \frac{p}{p^0} \right)$$

Where

$y_i$  is the mole fraction of species  $i$

$b_{ij}$  is the stoichiometric coefficient of component  $j$  in species  $i$

$P^0$  is the standard pressure (1 bar)



## Random Substitutional Solutions

- There is random substitution of one type of constituents by other. Ex.: molten steel,  $\alpha$ -brass
- In its most general form, molar Gibbs energy of a solution phase consists of five parts

$$G = G^{ref} + G^{conf} + G^E + G^{pres} + G^{phys}$$

- $G^E$  is the excess Gibbs energy term
- For an *ideal solution* last three terms are absent, i.e.

$$G = G^{ref} + G^{conf}$$
$$G = G^{ref} - TS^{conf}$$
$$G = \sum_i x_i (G_i^o - H_i^{SER}) + RT \sum_i x_i \ln x_i$$

- For real solutions,  $G^E \neq 0$

$$G = \underbrace{\sum_{i=A,B} x_i (G_i^o + H_i^{SER})}_{G^{\text{ref}}} + \underbrace{RT \sum_{i=A,B} x_i \ln x_i}_{G^{\text{conf}}} + \underbrace{x_A x_B L_{A,B}}_{G^E}$$

- Substitutional regular solution model for a phase in binary A-B

- Real solutions often need a model for composition dependence of  $G^E$  with many adjustable parameters
- Power series: polynomial functions of composition
- Examples: Redlich-Kister, Legendre, Margules, ...
- All polynomial models are inter-convertible
- Some are numerically superior than others
- Redlich-Kister (R-K) polynomial is the most commonly used since they have some advantages in multi-component systems

$$G^E = x_A x_B L_{A,B}(x_A, x_B)$$

$$G^E = x_A x_B \sum_{\nu=0}^n {}^{\nu}L_{A,B}(x_A - x_B)^{\nu}$$

- Where  ${}^{\nu}L_{A,B}$  are the R-K interaction parameters
- Their T dependence can be described using

$${}^iL_{A,B} = {}^i a + {}^i b T + {}^i c T \ln T$$



# Extrapolation models for ternary and higher order systems

- Allow us to express excess property of a multicomponent system in terms of excess properties of limiting binary systems.
- It is accomplished through geometrical extrapolation of binary  $G^E$ , resulting in

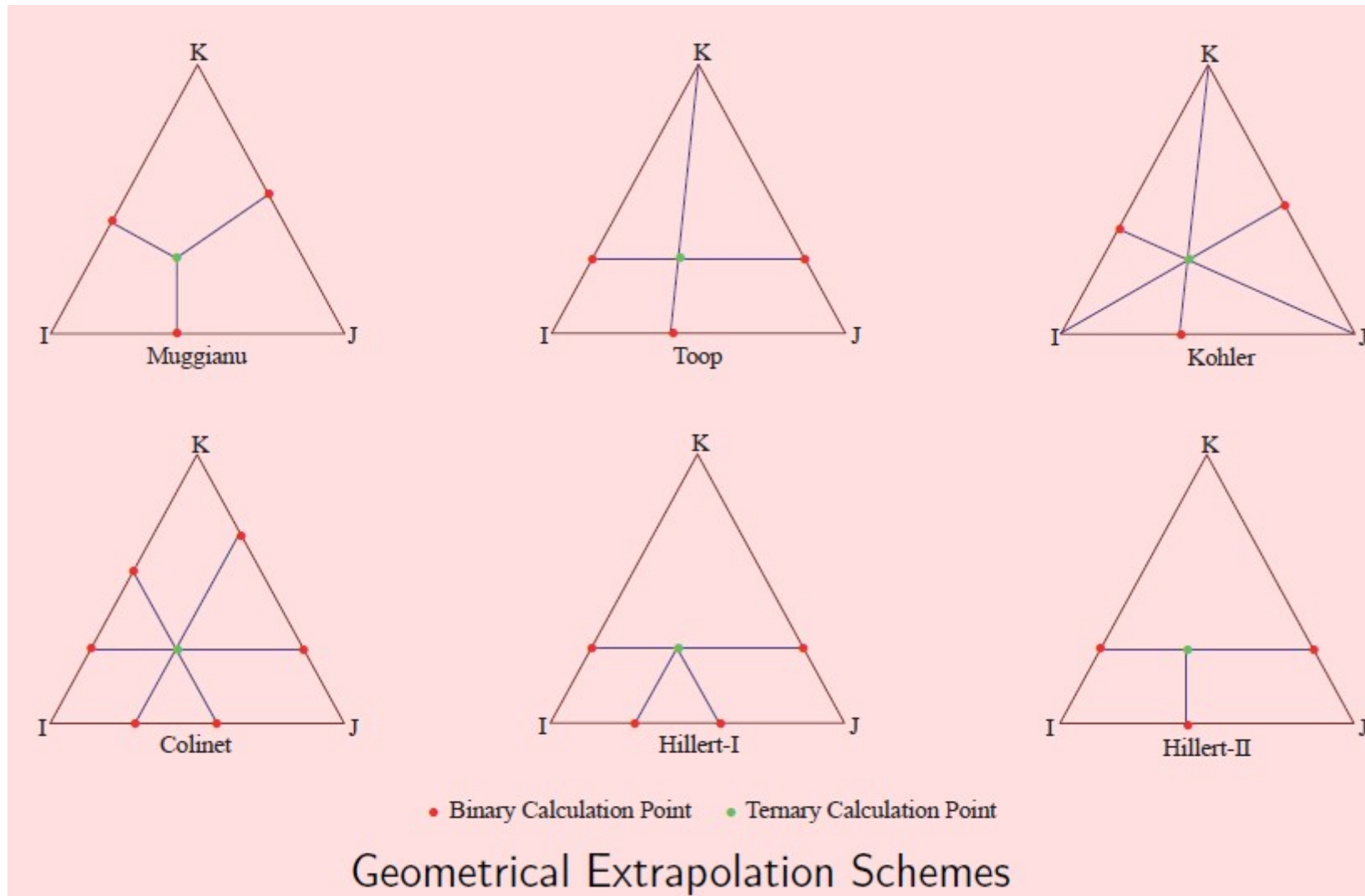
$$\left(G^E\right)_{m-c} = \left(G^E\right)_{expld} + \textit{Correction - terms}$$

- Correction terms are

$$\text{ternary} \rightarrow X_i X_j X_k L_{i,j,k}$$

$$\text{quaternary} \rightarrow X_i X_j X_k X_l L_{i,j,k,l}$$

- Many schemes exists for the extrapolation



- Muggianu extrapolation scheme:

$$\left(G^E\right)_{\text{expld}} = \sum_{i=1}^{c-1} \sum_{j=i+1}^c \frac{X_i X_j}{V_{ij} V_{ji}} \left(G^E(V_{ij}, V_{ji})\right)_{i-j}$$

- Where,  $\left(G^E(V_{ij}, V_{ji})\right)_{i-j} = V_{ij} V_{ji} L_{i,j}(V_{ij}, V_{ji})$

$$V_{ij} = X_i + \frac{(1 - X_i - X_j)}{2} \quad V_{ji} = X_j + \frac{(1 - X_i - X_j)}{2}$$

- The quantity  $\frac{X_i X_j}{V_{ij} V_{ji}}$  is a weighting factor
- If we use R-K polynomials for  $\left(G^E(V_{ij}, V_{ji})\right)_{i-j}$  above equation can be considerably simplified

$$\left(G^E\right)_{\text{expld}} = \sum_{i=1}^{c-1} \sum_{j=i+1}^c \left(G^E(V_{ij}, V_{ji})\right)_{i-j}$$

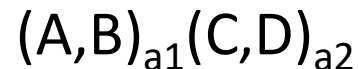
- In the case of ternary system R-K-M scheme gives

$$G^E = \sum_{i=1}^2 \sum_{j=i+1}^3 X_i X_j \sum_{\nu=0}^{n_{i-j}} {}^{\nu} L_{i,j} (X_i - X_j)^{\nu} + X_i X_j X_k L_{i,j,k}$$

# Sublattice Formalism

- General modeling concept for many types of phases: Stoichiometric compounds, random substitutional solutions, interstitial solutions, intermediate phases with homogeneity range, ionic and non-ionic melts, etc.
- Based on Temkin's model (1945) for molten salts  
Salt system –  $(\text{Cs}^+, \text{Na}^+)(\text{Cl}^-, \text{F}^-)$ 
  - Cations and anions occupy separate sublattices and mixing is random on each sublattice.
- Hillert and Staffanson (1970)

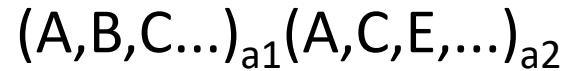
Formula unit containing two sublattices



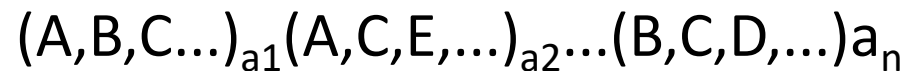
$a_s \rightarrow$  relative number of lattice sites on sublattice 's'

$$\sum_s a_s = 1 \quad (\text{or a positive integer})$$

- Harvig (1971)

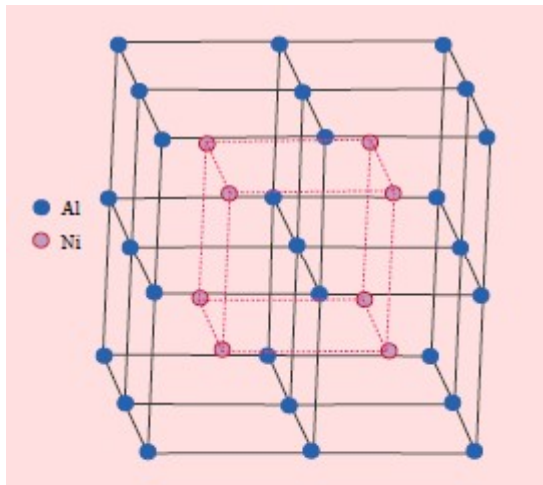


- Arbitrary number of constituent species
- Sundman and Ågren (1981): general formalism



- A,B, C,....=atoms, molecules, charges species, vacancy
- Arbitrary number of constituent species and sublattices allowed
- For solid phases: Compound Energy Formalism (CEF)
- A species is permitted only once in each sublattice
- Suitable for computer implementation

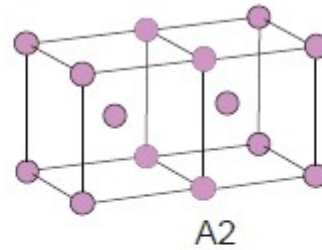
## Crystallographic data for the AlNi phase



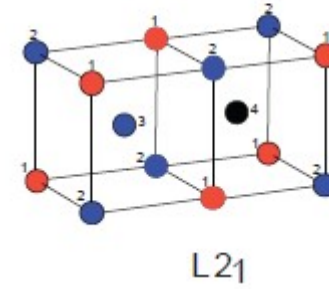
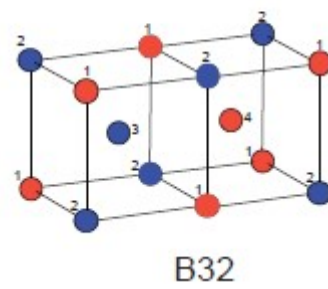
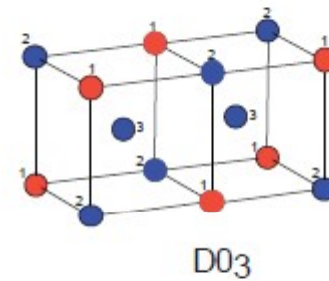
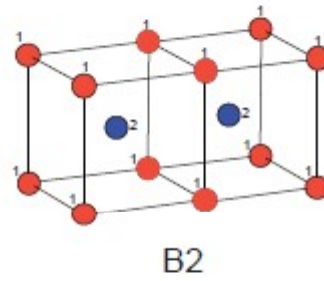
Prototype	CsCl	
Space Group	Pm/3m(221)	
Pearson Symbol	<i>cP</i> 2	
Strukturbericht	B2	
Point Group Symmetry	m/3m	m/3m
Site Occupation	Al	Ni
Wykoff Notation	1a	1b
Coordination Number	8	8

Sublattice model for perfectly ordered AlNi: (Al)<sub>1</sub>(Ni)<sub>1</sub>

Disordered bcc lattice



Ordered bcc lattice



● A or B   ● A   ● B   ● C

Figure 2.5: Sublattices of bcc based structures. (1, 2, 3 and 4 denotes sublattices)

- Site fractions: mole fractions defined for a particular sublattice.

$${}^s y_i = \frac{{}^s N_i}{{}^s N}$$

- ${}^s y_i$  → site fraction of constituent species  $i$  on sublattice  $s$
- ${}^s N_i$  → number of sites occupied by species  $i$  in sublattice  $s$
- ${}^s N$  → total number of sites on sublattice  $s$ , given by

$${}^s N = {}^s N_{va} + \sum_i {}^s N_i$$

- ${}^s N_{va}$  → number of vacant sites on sublattice  $s$
- Sum of site fractions in a given sublattice is unity

$${}^s Y_{va} + \sum_i {}^s y_i = 1$$



- Mole fraction of a component j can always be calculated from site fractions of constituent species

$$X_j = \sum_i \left( \frac{{}^i b_j}{\sum_k {}^i b_k} \left( \frac{\sum_s a_s {}^s y_i}{\sum_s a_s (1 - {}^s y_{va})} \right) \right)$$

- ${}^i b_j \rightarrow$  stoichiometry of component j in species i.  
e.g. (A, AB)<sub>1</sub>(B, Va)<sub>3</sub>

$$X_B = \frac{0.5 {}^1 y_{AB} + 3 {}^2 y_B}{1 + 3(1 - {}^1 y_{Va})}$$

- If all constituent species are monoatomic

$$X_j = \frac{\sum_s a_s {}^s y_j}{\sum_s a_s (1 - {}^s y_{va})}$$

- Site fractions have mixed character of an external and internal state variable
- They are obtained through Gibbs energy minimization
  - In simple cases site fractions can be calculated from mole fractions without resorting to a Gibbs energy minimization
  - For this, the number of independent site fractions should be equal to the number of independent mole fractions

e.g.  $(A,B)_{a_1}(C,Va)_{a_2}$

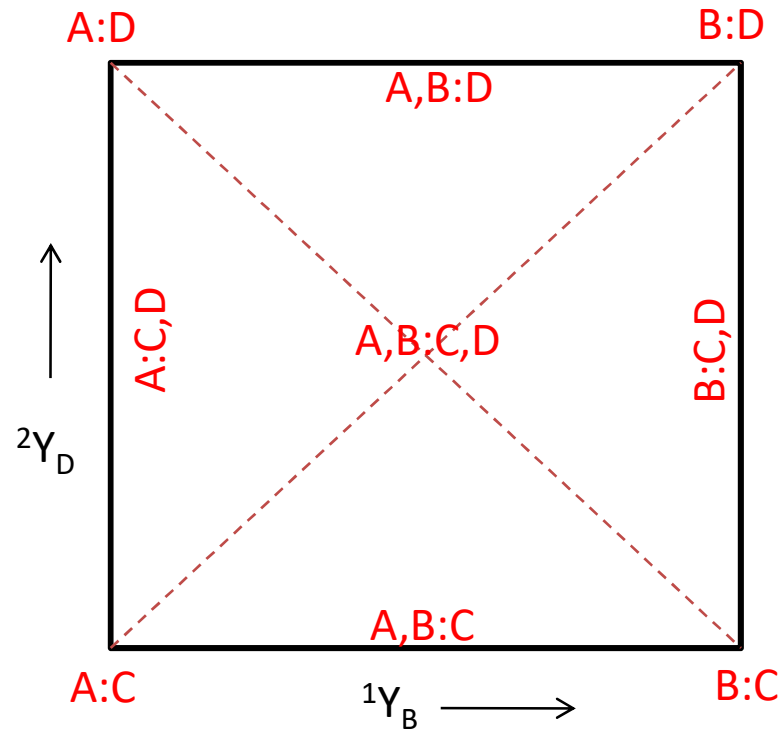
$$X_A = \frac{a_1^1 y_A}{a_1 + a_2(1 - {}^2 y_{Va})}$$

$$X_B = \frac{a_1^1 y_B}{a_1 + a_2(1 - {}^2 y_{Va})}$$

$${}^1 y_A = \frac{X_A}{1 - X_C}$$

$${}^2 y_A = \left( \frac{a_1}{a_2} \right) \frac{X_C}{1 - X_C}$$

- Constituent array: A unit of the sublattice model having one or more constituent species on each sublattice
- Constituent array of  $Z^{\text{th}}$  order:  $I(z)$
- Constituent arrays of  $(A,B)_{a_1}(C,D)_{a_2}$ 
  - $I(0) \rightarrow A:C, A:D, B:C$  and  $B:D$
  - No interacting constituent species on a sublattice
  - $I(1) \rightarrow A:(C,D), B:(C,D)$  and  $(A,B):C$   $(A,B):D$
  - One interacting constituent species on a sublattice
  - $I(2) \rightarrow (A,B):(C,D)$
  - Two interacting constituent species on two sublattices or three interacting constituent species on one sublattice.
- “.” is used to separate two sublattices and a “:” to separate two constituent species within a sublattice
- Generally, we need not consider constituent arrays larger than  $I(2)$ , since their influence on  $G$  is negligible
- $I(0)$  are also known as the “end members” of the sublattice model.



Composition domain for  $(A,B)_{a1}(C,D)_{a2}$

- $I(0)$  represent the corners of the composition domain of the model,  $I(1)$  the edges, and  $I(2)$  the diagonals

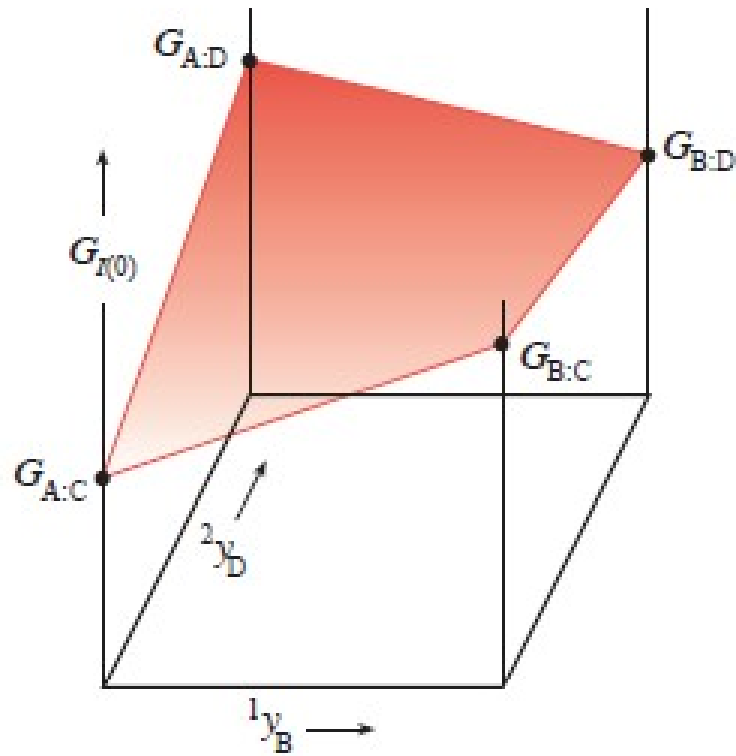


Fig.: Reference surface for  $(A,B)_{a_1}(C,D)_{a_2}$  at constant  $p$  and  $T$

- Reference surface is non-planar, defined by the function

$$G^{ref} = {}^1y_A ({}^2y_C G_{A:C}^o + {}^2y_D G_{A:D}^o) + {}^1y_B ({}^2y_C G_{B:C}^o + {}^2y_D G_{B:D}^o)$$

- It is planar when

$$G_{A:C}^o + G_{B:D}^o = G_{B:C}^o + G_{A:D}^o$$

- For the general case, we can write

$$G^{ref} = \sum_{I(0)} \left( \prod_{I(0)}^s y_i \right) G_{I(0)}^o$$

- $\prod_{I(0)}^s y_i$  is a product of site fractions from I(0)
- The degree of the product is equal to the number of sublattices having mixing of constituent species.

$$\sum_{I(0)} \prod_{I(0)}^s y_i = 1$$

- $G_{I(0)}^o$  is the Gibbs energy per mole of formula unit (mfu) of the end member compound –i.e. the compound energy
- Compound energy is the most significant part of the Gibbs energy of the phase and for this reason the sublattice model is popularly known as the “Compound Energy Formalism (CEF)”
- The number of terms (number of end members) is obtained by taking the product of number of species on each sublattice.
- $G_{I(0)}^o = f(T, p = 1bar)$ , may have several adjustable parameters.

- Constituent species on each sublattice mix randomly
- For the specific case of  $(A,B)_{a_1}(C,D)_{a_2}$  configurational entropy is given by

$$S^{conf} = -R \left[ a_1 \left( {}^1y_A \ln({}^1y_A) + {}^1y_B \ln({}^1y_B) \right) + a_2 \left( {}^2y_C \ln({}^2y_C) + {}^2y_D \ln({}^2y_D) \right) \right]$$

- For the general case,  $S^{conf}$  can be written as

$$S^{conf} = -R \sum_s a_s \sum_i {}^s y_i \ln({}^s y_i)$$

- The ideal Gibbs energy of mixing is given by

$$G^{conf} = -T.S^{conf} = -RT \sum_s a_s \sum_i {}^s y_i \ln({}^s y_i)$$

- Excess Gibbs energy ( $G^E$ ) represents interaction between constituent species within same sublattice
- $G^E$  involves only constituent arrays larger than  $I(0)$
- For the specific case of  $(A,B)_{a1}(C,D)_{a2}$ :

$$G^E = {}^1y_A {}^1y_B \left( {}^2y_C L_{A,B:C} + {}^2y_D L_{A,B:D} \right) +$$

$${}^2y_C {}^2y_D \left( {}^1y_A L_{A:C,D} + {}^1y_B L_{B:C,D} \right) +$$

$${}^1y_A {}^1y_B {}^2y_C {}^2y_D L_{A,B:C,D}$$

- $L_{I(Z)} \rightarrow$  interaction energies between constituent species in  $I(Z)$
- Influence of  $L_{I(1)}$  interaction terms ( $L_{A,B:C}$ ,  $L_{A,B:D}$ ,  $L_{A:C,D}$  and  $L_{B:C,D}$ ) are along the sides of the composition domain
- Influence of  $L_{I(2)}$  term ( $L_{A,B:C,D}$ ) is along the diagonals, maximum effect being at the center of the composition domain



- One may introduce further composition dependence for GE by expanding LI(1) terms as R-K polynomials for example,

$$L_{A,B:C} = \sum_{\nu=0}^k {}^{\nu}L_{A,B:C} \left( {}^1y_A - {}^1y_B \right)^{\nu}$$

- ${}^{\nu}L_{I(Z)}$  are adjustable model parameters
- They may have a temperature dependence

$${}^{\nu}L_{A,B:C} = {}^{\nu}a_{A,B:C} + {}^{\nu}b_{A,B:C}T + \dots$$

- Only in rare cases we need to expand  $L_{I(2)}$  terms

$$L_{A,B:C,D} = {}^0L_{A,B:C,D} + {}^1L_{A,B:C,D} \left( {}^1y_A - {}^1y_B \right) + {}^2L_{A,B:C,D} \left( {}^2y_C - {}^2y_D \right)$$